

The Crystal Structure of 2-Oxazolidinone

BY JUNE W. TURLEY

Chemical Physics Research Laboratory, The Dow Chemical Company, Midland, Michigan 48640, U.S.A.

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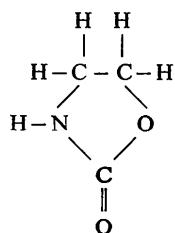
Crystals of 2-oxazolidinone, $C_3H_5NO_2$, are in monoclinic space group $P2_1/c$ with four molecules in a unit cell of dimensions $a = 7.313$ (4), $b = 5.672$ (5), $c = 9.970$ (9) Å and $\beta = 110.78$ (11)° measured at 24°C; $D_x = 1.50$ g.cm⁻³ and $D_m = 1.42$ (2) g.cm⁻³. Visual data, comprising 551 independent reflections above background, were collected with Cu $K\alpha$ radiation (1.5418 Å) and scaled from a Wilson plot. Iterative application of Sayre's equation in three dimensions led immediately to the solution of the structure. Refinement by full-matrix least-squares procedures gave a final residual index $R_1 = 0.132$ with hydrogen atoms included. The molecule is essentially planar with a root-mean-square deviation of 0.05 Å from the least-squares plane. There is one N-H···O hydrogen bond (2.858 Å) per molecule resulting in chains in the y direction, and no other close intermolecular distances. Bonds clockwise around the ring are C-C: 1.497 (10), C-O: 1.453 (8), O-C: 1.356 (10), C-N: 1.301 (8), and N-C: 1.466 (7) Å. The carbonyl C-O distance is 1.210 (8) Å. Angles in the ring are 100.3 (6), 106.0 (5), 108.6 (5), 110.2 (5), and 113.3 (6)° at C, C, O, C, and N respectively.

Introduction

Chemical studies (Walles, 1961) have shown that the ability of 2-oxazolidinone to form molecular complexes readily with materials such as phenols, aspirin, saccharin, iodine, etc., can be attributed almost entirely to the oxazolidinone ring. The present investigation was undertaken to elucidate the molecular geometry as a basis for studying this complexing ability. Structures based on the oxazole ring system have not been studied previously, although a recent report has been made (Amirthalingam & Muralidharan, 1969) on a thiazole compound said to be a potential radioprotective agent.

Experimental

Crystals of 2-oxazolidinone, prepared by W. E. Walles,



are small, colorless prisms elongated along a . One such crystal, approximately 0.2 mm on a side and mounted on the b axis, gave equi-inclination Weissenberg photographs showing C_{2h} diffraction symmetry and extinctions for $h0l$ with $l = 2n + 1$ and $0k0$ with $k = 2n + 1$, thereby establishing space group $P2_1/c$.

Lattice parameters, initially measured from Weissenberg single-crystal photographs taken with Cu $K\alpha$ (1.5418 Å) radiation, were subsequently refined using X-ray powder-diffraction data, taken at 24°C with crystal monochromatized Cu $K\alpha_1$ radiation (1.54050 Å)

on an AEG* Guinier camera using Seeman-Bohlin focusing (Hofman & Jagodzinski, 1955). A 7.5 μ Al foil sample support was used to provide reference lines. Three cycles with the method of axial ratios (Frevel, 1964) gave $a = 7.313 \pm 0.004$, $b = 5.672 \pm 0.005$, $c = 9.970 \pm 0.009$ Å, and $\beta = 110.78 \pm 0.11$ °. The crystal density measured by displacement in cyclohexane is 1.423 g.cm⁻³, in fair agreement with the calculated density, 1.495 g.cm⁻³, assuming four molecules per unit cell.

Intensity data for five levels along the b axis were read from equi-inclination Weissenberg multiple-film exposures made using Ni-filtered Cu $K\alpha$ (1.5418 Å) radiation. The films were read twice, independently by two different observers who compared them visually with a scale calibrated by a reflection from the same crystal. The data sets were then scaled internally using Wilson scale factors and were merged. The copper sphere allowed 883 reflections of which 551 were observed. The intensities were corrected for Lorentz and polarization effects, but no absorption correction was made ($\mu = 10.46$ cm⁻¹). Assuming $\sigma(F^2) = 0.10F^2$ for a given unique reflection, the expression $\sigma(F) = \sigma(F^2)/2F$ was used to estimate the σ 's needed for computing the weights $w = 4F^2/\sigma^2(F^2)$ used in the least-squares refinement (Gvildys, 1967), where the quantity $\sum w(F_o - F_c)^2$ was minimized. Reflections 011, $\bar{2}11$, and $\bar{2}12$ were excluded from least-squares refinements because of severe secondary extinction effects, seen as a factor-of-two difference between observed and calculated structure factors for these data.

Solution and refinement of the structure

Iterative application of Sayre's (1952) equation in three dimensions led immediately to the solution of the struc-

* Allgemeine Elektrizitäts-Gesellschaft.

ture. The equation is:

$$s(E_a) = s \left(\sum_{a=b+c} E_a E_c \right)$$

where s means 'sign of', a , b , c are the vectors ($hk\ell$) for the reflexions a , b , c , and E_a , E_b , & E_c are normalized structure factors for a , b , c . A set of 182 reflections having $|E| \geq 1.5$, and a starting set comprising three origin-determining reflections assigned positive phases and four additional reflections with both positive and negative phases allowed, were used with a completely automated computer program (Long, 1965). E maps (Gvildys, 1965a) were calculated for the two unique solutions found, and the structure was clearly shown

on one while the other gave a map with too few peaks and lower peak heights. An initial calculation with positions read from the E map gave a residual index $R = \sum |F_o| - |F_c| / \sum |F_o|$ of 0.334. Using a full-matrix least-squares program (Gvildys, 1967), refinement was carried out as described below. Four cycles of isotropic refinement defining all peaks as carbon atoms (Ibers, 1962) decreased R to 0.248. At this point, atoms were redefined as carbon, nitrogen, and oxygen based on refined thermal parameters and the known atomic sequence in the molecule. Next, three cycles of isotropic refinement, followed by three cycles of anisotropic refinement, gave $R = 0.155$, and a difference map showed positions for the five hydrogen atoms in the molecule.

Table 1. Final fractional coordinates and anisotropic thermal displacement ($\text{\AA} \times 10^3$)* for 2-oxazolidinone

Fractional coordinates of non-hydrogen atoms are multiplied by 10^4 and those of the hydrogen atoms by 10^3 . E.s.d.'s based on the least-squares refinement are given in parentheses.

	x	y	z	$U_{11}^{1/2}$	$U_{22}^{1/2}$	$U_{33}^{1/2}$	$U_{12}^{1/2}$	$U_{13}^{1/2}$	$U_{23}^{1/2}$
C(1)	2507 (12)	-3051 (13)	5530 (7)	254 (70)	153 (72)	196 (60)	35 (58)	190 (57)	100 (52)
C(2)	3619 (11)	-1133 (12)	6515 (7)	237 (66)	185 (68)	176 (56)	55 (56)	89 (54)	67 (51)
C(3)	2047 (9)	502 (13)	4312 (6)	181 (56)	191 (32)	162 (53)	47 (52)	75 (49)	25 (49)
N	1433 (9)	-1664 (12)	4251 (6)	234 (60)	181 (69)	171 (53)	34 (54)	77 (50)	25 (47)
O(1)	3406 (7)	952 (8)	5624 (4)	207 (50)	188 (54)	159 (45)	88 (44)	60 (42)	53 (41)
O(2)	1524 (8)	2034 (9)	3411 (5)	237 (55)	242 (62)	186 (49)	37 (49)	44 (46)	138 (47)
H(1)	364 (12)	-389 (13)	534 (8)	220	220	220	—	133	—
H(2)	191 (12)	-405 (15)	571 (9)	222	222	222	—	134	—
H(3)	507 (15)	-139 (17)	691 (11)	286	286	286	—	172	—
H(4)	313 (14)	-64 (15)	729 (10)	270	270	270	—	163	—
H(5)	95 (13)	-241 (16)	357 (10)	244	244	244	—	147	—

* These values are the square roots of the mean-square amplitude tensors, U_{ij} , based on the expression $\exp(-h^2\beta_{11} + k^2\beta_{22} + l^2\beta_{33} + 2hk\beta_{12} + 2hl\beta_{13} + 2kl\beta_{23})$ and relations such as $\beta_{12} = 2\pi^2 U_{12} a^* b^*$ (Cruickshank, 1965).

Table 2. Observed and calculated structure factors for 2-oxazolidinone

H	L	FORS	FCALC	H	L	FORS	FCALC	H	L	FORS	FCALC	H	L	FORS	FCALC	H	L	FORS	FCALC	H	L	FORS	FCALC							
*** K = 0 ***				# 12	887	926	3 1	287	273	# 1	171	127	0 8	0	137	-1 7	704	569	-5 7	324	268	4 5	514	375	-5 9	250	148	3 6	227	235
				# 5	541	555	2 3	1217	975	# 4	167	1093	0 6	0	592	-1 6	832	590	-1 6	224	174	4 6	511	369	-1 10	423	242	3 7	331	342
				# 5	198	205	3 5	121	607	# 4	236	211	0 11	0	267	-1 6	758	758	-1 6	188	142	5 7	508	358	-1 12	423	242	3 7	331	342
				# 5	198	205	3 5	121	607	# 4	236	211	0 11	0	267	-1 6	758	758	-1 6	188	142	5 7	508	358	-1 12	423	242	3 7	331	342
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				# 5	198	205	3 5	121	607	# 4	236	211	0 11	0	267	-1 6	758	758	-1 6	188	142	5 7	508	358	-1 12	423	242	3 7	331	342
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Finally, three cycles with all atoms anisotropic except the hydrogen (Stewart, Davidson & Simpson, 1965) gave an *R* value of 0.132. Reasonable convergence was indicated by a maximum coordinate shift of 0.04σ for non-hydrogen atoms and 0.17σ for hydrogen atoms, and a value of 1.97 for the indicator $[\sum(w\Delta F^2)/(NO-NV)]^\frac{1}{2}$, where *NO* is the number of observations and *NV* is the number of parameters varied. The refined atomic parameters are listed in Table 1, and the observed and calculated structure factors are compared in Table 2.

Discussion

Fig. 1 shows the molecular structure (Johnson, 1965). The molecule is approximately planar with a root-mean-square deviation of 0.05 Å from the least-squares plane of all atoms (Gvildys, 1965b), plane 1 in Table 3, from which atoms C(1) and C(2) show the largest deviation. Plane 2 is fitted to atoms C(3), O(1), N, and O(2), and plane 3 to atoms C(1), C(2), O(1), and N; the dihedral angle between them is $2^\circ 33'$, showing a significant deviation from planarity when compared with the estimated errors of 0.6° in the bond angles.

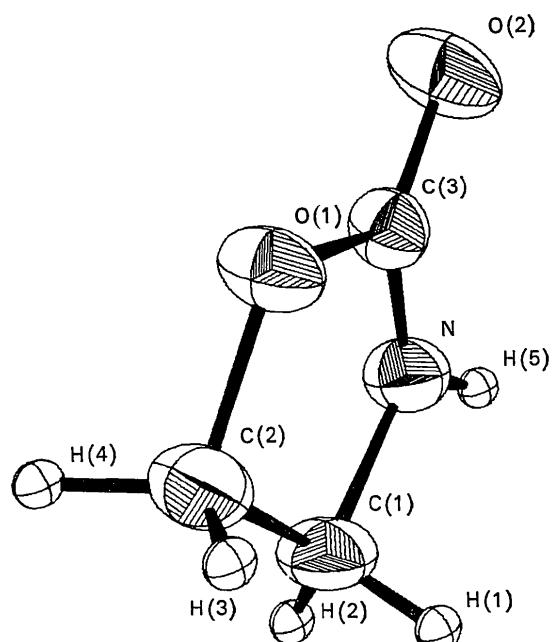
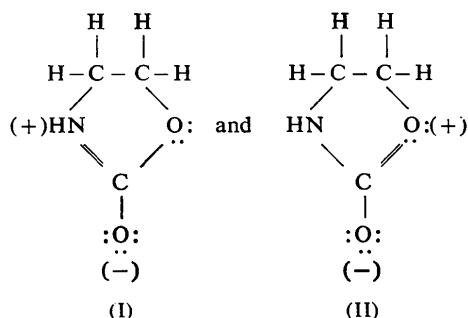


Fig. 1. Molecular structure of 2-oxazolidinone with atoms represented as 50% probability thermal-motion ellipsoids.

Although the bond distances and angles, given in Table 4 (Gvildys, 1964), show some differences from expected values (Sutton, 1965) they agree well with values found in other lactone structures (Koenig, Chiu, Krebs & Walter, 1969; Kalyani & Vijayan, 1969). The C(1)-C(2) distance is 0.04 Å less than the expected value of 1.537 Å for a simple C-C bond, C(2)-O(1) is 0.02 Å longer than the expected C-O single bond length of 1.435 Å, and C(3)-O(1) is 0.04 Å longer than the analogous C-O bond (1.312 Å) in organic acids but much shorter than a C-O single bond distance. Also, the C(3)-N distance of 1.301 Å is much shorter than expected (1.333 Å) for an N-C=O group, and the C(1)-N distance of 1.466 Å is shorter than the value of 1.472 Å usually observed for trivalent nitrogen. Although these observed bond distances have moderately high estimated errors, the differences discussed above are large enough to suggest that the major resonance forms for the structure are:



Here, the N=C-O portion of the molecule can act as a dipole on a polarizable molecule, suggesting an explanation for the unusual complexing ability of this ring system.

The molecular packing is shown in Fig. 2. Hydrogen-bonded chains are formed in the *y* direction by the 2.862 Å N-H...O(2) interaction, resulting in planar strips slightly more than two molecules wide and aligned to suggest some dipole-dipole interaction between them in the *x* direction. Only van der Waals approaches of CH₃ groups occur in the *z* direction.

The author wishes to express her appreciation to D. R. Petersen who collected the data set, F. P. Boer who was particularly helpful in preparing the Figures and encouraging the work, and B. K. Moss who processed much of the computing.

Table 3. *Molecular planes in 2-oxazolidinone*

Equation of the plane is $M_1x + M_2y + M_3z = D$, where *D* is the origin-to-plane distance and *x*, *y*, *z* are fractional coordinates in the crystal system. The Δ 's are the distances of the atoms from the fitted plane.

Plane	<i>M</i> ₁	<i>M</i> ₂	<i>M</i> ₃	<i>D</i>	$\Delta C(1)$	$\Delta C(2)$	$\Delta O(1)$	$\Delta C(3)$	ΔN	$\Delta O(2)$
1	6.687	-1.428	-6.189	-1.381	0.071 Å	-0.069 Å	0.042 Å	0.010 Å	-0.054 Å	-0.001 Å
2	6.554	-1.606	-6.354	-1.490	0.110†	-0.095†	-0.003	0.012	-0.004	-0.005
3	6.684	-1.406	-6.228	-1.409	0.069	-0.071	0.049	0.021†	-0.047	0.017†

† These atoms were not included when fitting the plane.

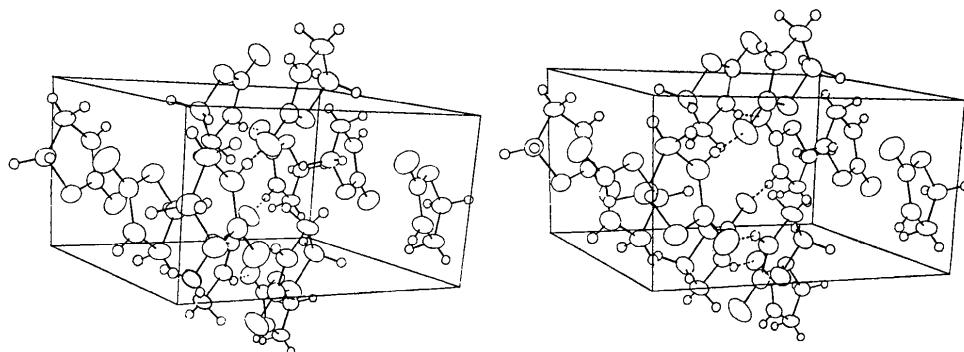


Fig. 2. Stereo view of molecular packing in 2-oxazolidinone with one unit cell outlined.

Table 4. Interatomic distances and angles in 2-oxazolidinone

Errors given in parentheses are in the last one or two significant figures.

C(1)–C(2)	1.497 (10) Å	N—C(1)–C(2)	100.3 (6)°
C(2)–O(1)	1.453 (8)	H(1)–C(1)–N	108.0 (43)
O(1)–C(3)	1.356 (10)	H(1)–C(1)–C(2)	99.9 (43)
C(3)–N	1.301 (8)	H(1)–C(1)–H(2)	104.8 (82)
N—C(1)	1.466 (7)	H(2)–C(1)–N	115.2 (61)
C(3)–O(2)	1.210 (8)	H(2)–C(1)–C(2)	126.7 (61)
C(1)–H(1)	1.03 (8)	C(1)–C(2)–O(1)	106.0 (5)
C(1)–H(2)	0.77 (8)	H(3)–C(2)–C(1)	113.8 (58)
C(2)–H(3)	1.00 (10)	H(3)–C(2)–O(1)	103.1 (56)
C(2)–H(4)	1.00 (9)	H(3)–C(2)–H(4)	111.0 (79)
N—H(5)	0.77 (9)	H(4)–C(2)–O(1)	104.8 (50)
N.....O(2)	2.862 (7)	H(4)–C(2)–C(1)	116.6 (53)
H(5)...O(2)	2.18 (9)	C(2)–O(1)–C(3)	108.6 (5)
		O(1)–C(3)–N	110.2 (5)
		O(1)–C(3)–O(2)	120.7 (6)
		N—C(3)–O(2)	129.0 (6)
		C(3)–N—C(1)	113.3 (6)
		H(5)–N—C(1)	114.2 (70)
		H(5)–N—C(3)	127.0 (70)

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